

Isotopic composition of CO₂ in the coma of 67P/Churyumov-Gerasimenko measured with ROSINA/DFMS

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ABSTRACT

Context. Measurements of isotopic abundances in cometary ices are key to understanding and reconstructing the history and origin of material in the solar system. Comets are considered the most pristine material in the solar system. Isotopic fractionation (enrichment of an isotope in a molecule compared to the initial abundance) is sensitive to environmental conditions at the time of comet formation. Therefore, measurements of cometary isotope ratios can provide information on the composition, density, temperature, and radiation during formation of the molecules, during the chemical evolution from the presolar cloud to the protosolar nebula, and the protoplanetary disk before accretion in solid bodies. Most isotopic abundances of ¹²C/¹³C and ¹⁶O/¹⁸O in comets to date are in agreement with terrestrial abundances. Prior to the Rosetta mission, measurements of ¹²C/¹³C in comets were only available for HCN, CN, and C₂ and for ¹⁶O/¹⁸O in H₂O. Measurements of ¹²C/¹³C in comets were only available from ground based observations and remote sensing, while ¹⁶O/¹⁸O in H₂O had also been measured in-situ. To date, no measurements of the CO₂ isotopologues in comets were available.

Aims. This paper presents the first measurements of the CO₂ isotopologues in the coma of 67P/Churyumov-Gerasimenko (67P).

Methods. We analyzed measurements taken by the Double Focusing Mass Spectrometer (DFMS) of the ROSINA experiment on board the ESA spacecraft Rosetta in the coma of 67P.

Results. The CO₂ isotopologues results for 67P are: ¹²C/¹³C = 84 ± 4, ¹⁶O/¹⁸O = 494 ± 8, and ¹³C¹⁶O₂/¹²C¹⁸O¹⁶O = 5.87 ± 0.07. The oxygen isotopic ratio is within error bars compatible with terrestrial abundances but not with solar wind measurements.

Conclusions. The carbon isotopic ratio and the combined carbon and oxygen isotopic ratio are slightly (14%) enriched in ¹³C, within 1σ uncertainty, compared to solar wind abundances and solar abundances. The small fractionation of ¹²C/¹³C in CO₂ is probably compatible with an origin of the material in comets from the native cloud.

Key words. comets: individual: 67P – astrochemistry – methods: data analysis – Kuiper belt objects: individual: 67P – comets: general

1. Introduction

To understand and reconstruct the history and origin of material in the solar system, measurements of isotopic abundances are key. Since comets are considered to contain some of the most pristine material in the solar system, measurements of isotopic abundance ratios of different atoms in different molecules are important. Isotopic fractionation describes the processes which affect the relative abundances of the isotopes. Isotopic fractionation is sensitive to environmental conditions at the time of comet formation. Therefore, measurements of isotope ratios provide information on the composition, density, temperature, and radiation during formation of the molecules and during the chemical evolution from the presolar cloud to the protosolar nebula and the protoplanetary disk before accretion in solid bodies. This paper focuses on the isotopic abundance of ¹³C and ¹⁸O in CO₂

for 67P/Churyumov-Gerasimenko (hereafter: 67P). It is the first time that these isotopologues have been measured in a comet.

The local interstellar medium (ISM) has an isotopic abundance ratio of ¹²C/¹³C (CN, CO and H₂CO) of 68 ± 15 and this ratio changes over galactic timescales as a function of galactocentric distance and time (Milam et al. 2005). The modeling results of Roueff et al. (2015) for the prestellar core and a moderately dense cloud show that the ¹²C/¹³C ratios for various molecules in the ISM are highly sensitive to the chemical evolution time of the prestellar core or moderate dense cloud. Various chemical reactions are responsible for incorporating ¹³C in molecules and the transition from gas-phase atomic carbon towards CO controls the ¹³C enrichment. As long as there is still a relatively high carbon-atom concentration in the gas phase, there is enough free ¹³C to allow strong enrichment of CN through the ¹³C + CN reaction (Roueff et al. 2015). Because eventually

most of the carbon is locked into CO and the $^{12}\text{C}/^{13}\text{C}$ ratio is very close to the atomic $^{12}\text{C}/^{13}\text{C}$ ratio after 10^6 yr, this ratio will change very little. Similar ratios are incorporated in CO_2 through the reaction of CO with OH (Woods & Willacy 2009).

The present-day terrestrial isotopic abundance ratio of $^{12}\text{C}/^{13}\text{C}$ is 89 (Meija et al. 2016). Meija et al. (2016) list the isotopic composition of the elements, while the best measurement of isotopic abundances from a single terrestrial source (NBS 19) is cited as a measurement by Zhang et al. (1990). The terrestrial abundances by Lodders (2003) are based on Rosman & Taylor (1998) and use the same reference as standard for terrestrial $^{12}\text{C}/^{13}\text{C}$ abundances. However, there exists a lower ratio of 80 ± 1 for the solar photosphere (Ayres et al. 2006). The $^{12}\text{C}/^{13}\text{C}$ ratio of 68 for the local interstellar medium is considerably lower than terrestrial. An even higher ratio than the terrestrial ratio was found in the solar wind by Hashizume et al. (2004). This may be attributed to ^{12}C enrichment caused by the proximity of the Sun to one or more massive stars, which produce ^{12}C in their interiors (Timmes et al. 1995). These stars may have contaminated the solar protoplanetary nebula with ^{12}C -rich material during its formation (Woods & Willacy 2009). The solar wind measurements by Hashizume et al. (2004) are much older than the measurements of the oxygen isotopes in the solar wind by NASA's Genesis mission (McKeegan et al. 2011). According to the findings by McKeegan et al. (2011) the solar wind measurements do not directly represent the solar abundance. As the most plausible composition of the Sun, the intersection of the calcium-aluminium-rich inclusions of chondritic meteorites and the solar wind measurements was used. However, there are no measurements available from the Genesis mission for carbon isotopes and in Hashizume et al. (2004) no such calculations are provided. Therefore the measurements by Hashizume et al. (2004) are used as the $^{12}\text{C}/^{13}\text{C}$ solar abundance, although based on the findings by McKeegan et al. (2011) the real solar abundance might be somewhere between the solar wind measurements of 98 ± 2 and the value of 89 by Anders & Grevesse (1989).

It is important to note, however, that the carbon isotopic ratio is not constant within the solar system. While meteorites in general have isotopic ratios close to 89, interplanetary dust particles (e.g., Zinner et al. 1995: $^{12}\text{C}/^{13}\text{C}$ 7 to 4500) and cometary dust (e.g., Jessberger 1999: $^{12}\text{C}/^{13}\text{C} \sim 5000$) show deviations from this mean value. These solar system components conserved part of the original traces of nucleosynthesis in the star from which they obtained their carbon. Carbon trapped within the 50 nm thick surface layer of the lunar regolith exhibits a depletion of ^{13}C in solar wind of at least 10% relative to terrestrial and meteoritic carbon (Hashizume et al. 2004). Hashizume et al. (2004) provide a solar wind value for the $^{12}\text{C}/^{13}\text{C}$ ratio of 98 ± 2 . In primitive ice, by far the most abundant carbon-bearing molecule is CO, and CO_2 is chemically derived from CO. In the native molecular cloud, little ^{13}C enrichment or depletion is expected for these molecules. According to Woods & Willacy (2009), there might be carbon fractionation in the protoplanetary disk depending on the chemical reactions, species, and temperature along the direction perpendicular to the disk, and depending on the distance from the Sun in the case of a solar nebula-like disk. According to their model, C_2 would be depleted in ^{13}C compared to CN and HCN. Their model also predicts that CO_2 and CO would be enriched in ^{13}C compared to the solar value depending on the location in the disk. Among the molecules remotely observed in several comets are C_2 , CN, and HCN. Remote sensing measurements of a variety of comets show that there are no large differences in the carbon isotopic abundance depending on the molecule (see Table 3 and Fig. 1). Most measurements

Table 1. Solar abundances for carbon and oxygen isotopes.

Isotope	Solar abundance ^a	Terrestrial abundance ^b
$^{12}\text{C}/^{13}\text{C}$	98 ± 2^c	89
$^{16}\text{O}/^{17}\text{O}$	2798	2682
$^{16}\text{O}/^{18}\text{O}$	530	499
$^{12}\text{C}^{18}\text{O}/^{16}\text{O}$		
$^{13}\text{C}^{16}\text{O}_2$		2.79

Notes. Inverse values are provided to assist in comparing with other measurements. ^(a) McKeegan et al. (2011) no uncertainties provided. ^(b) Meija et al. (2016) based on Zhang et al. (1990) no uncertainties provided. ^(c) Solar wind measurements by Hashizume et al. (2004).

are compatible with solar or terrestrial abundances within error bars. Small variations between the different molecules are observed between different comets, but no significant variations are seen for a single species. The error bars of all the remote sensing measurements are more than 10% with the exception of the $^{12}\text{CN}/^{13}\text{CN}$ ratio of Manfroid et al. (2009), which represents an average over 21 comets. While on average C_2 and CN agree well within terrestrial abundances of 89, HCN appears to be somewhat depleted in ^{13}C although the deviation is less than 2σ . No remote sensing measurements for the isotopic abundance of $^{12}\text{C}/^{13}\text{C}$ in CO or CO_2 exist because, to this date, no pure rotation spectra in the infrared or microwave regions for CO_2 are available. Therefore there is a lack of detections of CO and CO_2 detections of rare isotopologues. The terrestrial $^{16}\text{O}/^{18}\text{O}$ ratio derived from measurements of Vienna Mean Standard Ocean Water is known to be 499 (Meija et al. 2016; Lodders 2003; and Anders & Grevesse 1989), while the ratio in solar wind samples captured by the Genesis mission provides a solar abundance of 530 (McKeegan et al. 2011). By mass balance the latter could be considered the abundance of the protoplanetary nebula. The measurements presented here are discussed in the context of both of these ratios. Remote sensing and in situ measurements for comets of oxygen isotopes are only available for H_2O and these measurements have rather large uncertainties. Most measurements are within error bars compatible with both known values. Solar and terrestrial isotopic abundances for $^{12}\text{C}/^{13}\text{C}$ and $^{16}\text{O}/^{18}\text{O}$ from the literature are given in Table 1. The purpose of this paper is to report the first in situ measurement of the isotopic abundance of carbon and oxygen in CO_2 for ^{13}C and ^{18}O . CO_2 is originally derived from CO (due to chemical reactions) and together they represent the largest carbon reservoir in the solar system (Lewis & Prinn 1980). Even if isotopic fractionation plays a role for minor carbon bearing species, it would be difficult to significantly change the $^{12}\text{C}/^{13}\text{C}$ ratio in CO and CO_2 (see Woods & Willacy 2009). The $^{12}\text{C}/^{13}\text{C}$ ratio in CO and CO_2 in comets may therefore provide an additional clue as to the initial isotopic ratio for C in the protoplanetary nebula. The Rosetta mission spent months exploring the coma of comet 67P/Churyumov-Gerasimenko, measuring this coma with a wide variety of instruments. One of the instruments on board is ROSINA (Rosetta Orbiter Spectrometer for Ion and Neutral Analysis Balsiger et al. 2007). ROSINA consists of three sensors, out of which the instrument of choice for isotopologues is the Double Focusing Mass Spectrometer (DFMS).

1.1. ROSINA DFMS Sensor on Rosetta

ROSINA/DFMS is a high-resolution mass spectrometer designed to resolve molecules of nearly identical mass and

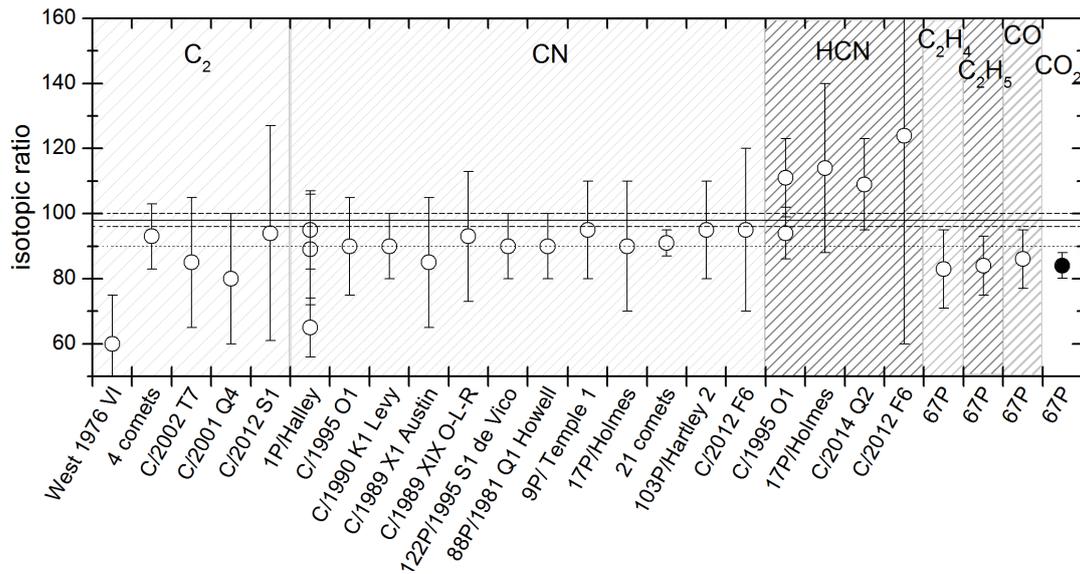


Fig. 1. Measurements of $^{12}\text{C}/^{13}\text{C}$ in different molecules for several comets (for references see Table 3). Several measurements of $^{12}\text{C}/^{13}\text{C}$ in the coma of 67P were taken by ROSINA/DFMS. The coma of 67P shows an enrichment of ^{13}C in CO_2 (black circle, this work) compared to solar abundance. For discussion on uncertainties see Sects. 4 and 5. The dotted horizontal line represents the terrestrial $^{12}\text{C}/^{13}\text{C}$ ratio. The solid horizontal line (dashed lines represent uncertainties) represents the solar wind $^{12}\text{C}/^{13}\text{C}$ ratio.

measure isotopic abundances. A detailed description of the instrument and its capabilities can be found in Balsiger et al. (2007), while data analysis methods are extensively described in several ROSINA/DFMS papers (e.g., Le Roy et al. 2015; and Calmonte et al. 2016). DFMS is a double focusing mass spectrometer with a high mass resolution and very high dynamic range, which separates species according to their mass to charge ratio (m/z). The high mass resolution of 3000 at 1% peak level allows the separation of, for example, ^{12}CH from ^{13}C , while the very high dynamic range of $\sim 10^{10}$ was very important for the encounter with 67P because the difference in densities from CO_2 (and H_2O) to the rarest of volatiles was expected to be four orders of magnitude, while the variations in H_2O and CO_2 outgassing were expected to present an additional three orders of magnitude difference from initial encounter to perihelion. The principles by which DFMS operates are as follows: cometary neutrals (both atoms and molecules) enter ROSINA/DFMS and are ionized by electron ionization. The formed ions are guided through the electrostatic analyzer and the magnet where they are separated according to their mass to charge ratio (m/z). Only ions with a specific m/z value have a stable trajectory through the instrument and DFMS sequentially measures masses between 13 to more than 100 u/e. One m/z scan takes approximately 20 s so that the entire mass range is typically covered in about 45 min. Particles with the selected m/z then hit the micro channel plate (MCP), where they release an avalanche of electrons. A linear detector array (LEDA) consisting of two individual rows with 512 pixels each then detects these electrons. The avalanche of electrons released, thus the gain of the detector, is controlled by the voltage applied over the MCP. The 16 separate voltage settings over the MCP are also called gain steps. The LEDA detects the electron charge of each individual pixel. The collimated ion beam typically hits the center of the MCP. This causes a preferential degradation of the center pixels of the MCP compared to the edge pixels of the MCP as a function of time. To account for this MCP ageing over the mission duration, the ion beam is regularly swept across all MCP/LEDA pixels and the results allow a correction factor for each individual pixel to be established.

DFMS measurements have already contributed significantly to cometary science by confirming and discovering cometary species (e.g., N_2 by Rubin et al. 2015; O_2 by Bieler et al. 2015), by providing a link between remote sensing and in situ coma measurements (Le Roy et al. 2015), by discovering high variations in local compositional densities (Hässig et al. 2015; and Luspay-Kuti et al. 2015), and by obtaining measurements of isotopic species with very low abundances (D/H by Altwegg et al. 2015).

1.2. Methodology for determining isotopic abundances with ROSINA/DFMS

Measurements on different gain steps, pixel gain across the MCP, and the time required to measure two individual isotopologue provide the largest uncertainties in measuring isotopic ratios with DFMS. Ideally, these uncertainties are minimized by comparing two mass peaks taken with the same gain step at the same location on the MCP and within as short amount of time as possible. These conditions are not always met and minimizing the uncertainties drives the data analysis. Three possibilities exist to measure isotopic ratios with DFMS: the first and second using two different masses, and the third using nearly the same mass (e.g., HD^{16}O and H_2^{17}O):

1. The species for an isotopic ratio measurement are on different masses. If the isotopic abundance ratio is within dynamic range of the LEDA and no other high abundance species is determining the gain of the detector, then the measurements are taken on the same gain step and only corrections for the pixel gain have to be taken into account. For measurements of species with the peak center on the same location of the detector, (which is normally the case for isotopologues) the correction due to the individual pixel gain is very small and measurements are therefore considered very precise.
2. The species for an isotopic ratio measurement are on different masses. If the species do not have a similar abundance and the peak center is shifted by several pixels, corrections for the different gain steps and the individual pixel gain have to be taken into account.

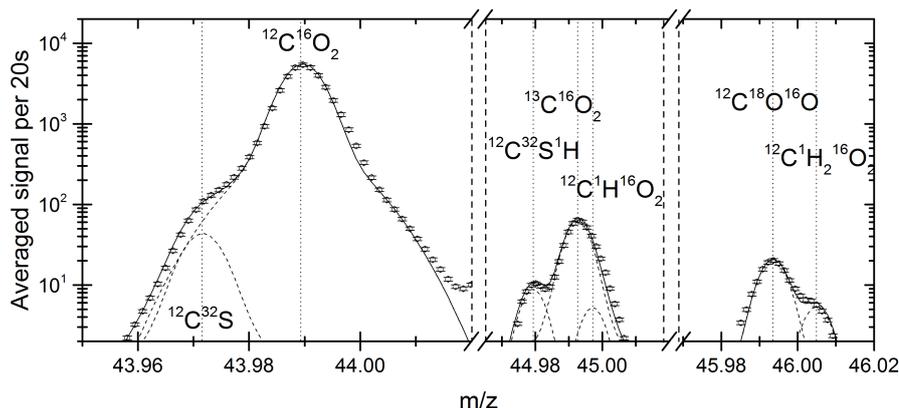


Fig. 2. Measured signal averaged over 87 co-added 20s spectra and detected species on m/z 44 u/e, 45 u/e, and 46 u/e. The circles show the measured points on the detector, while the dashed lines are the fit of the according peak fits. The solid line represents the sum of all fits. Error bars represent the uncertainty due to counting statistics (\sqrt{n}).

3. The isotopic ratio is measured on the same mass but a combination of several isotopes. Therefore, no gain step uncertainties have to be taken into account. However, due to the shift in position between the two peaks, a pixel gain correction is applied to account for the ageing of the multi channel plate (MCP) due to impinging ions on the same location of the MCP. This method is precise, but possible only for a few species, for example, ^{18}OH , HD^{16}O and H_2^{17}O , all on mass 19.

The first method is applied for the measurement of the CO_2 isotopes. On m/z 44, 45, 46 u/e the following CO_2 isotopologues are found:

44 u/e: $^{12}\text{C}^{16}\text{O}_2$;
 45 u/e: $^{13}\text{C}^{16}\text{O}_2$ and $^{12}\text{C}^{17}\text{O}^{16}\text{O}$;
 46 u/e: $^{12}\text{C}^{18}\text{O}^{16}\text{O}$ and $^{13}\text{C}^{17}\text{O}^{16}\text{O}$ and $^{12}\text{C}^{17}\text{O}_2$.

The abundance of molecules containing ^{17}O is one magnitude smaller than that for ^{13}C or ^{18}O and was therefore negligible.

A second source of uncertainty is the subtraction of neighboring peaks that overlap with isotopologues of interest.

In addition to the above-mentioned molecules, the following additional molecules are found on m/z 44, 45, 46 u/e, overlapping with the signal peak of CO_2 isotopes (see Fig. 2):

44 u/e: CS , $\text{C}_2\text{H}_4\text{O}$;
 45 u/e: CHS , CHO_2 ;
 46 u/e: CH_2S , NO_2 , CH_2O_2 .

The contribution of $\text{C}_2\text{H}_4\text{O}$ to the signal for CO_2 is clearly negligible, as it is at least three orders of magnitude lower. The difference in mass per charge between NO_2 and $^{12}\text{C}^{18}\text{O}^{16}\text{O}$ is only 0.002 u/e and cannot be resolved with DFMS. The abundance of NO_2 was therefore assumed to be negligible (see Sect. 4.), resulting in a possible overestimation of the signal abundance of $^{12}\text{C}^{18}\text{O}^{16}\text{O}$. As shown later, there is good observational data to substantiate this assumption.

The peak shape of DFMS is well known and signals of overlapping peaks are resolved by fitting a double Gaussian to the measured points. The peak location is known and the width of the peak is the same for all species on an integer mass spectrum, which then leaves the amplitudes for the different peaks and the common FWHM (Full Width Half Maximum) as fit parameters. This peak-fitting method was already successfully used to separate H_2^{17}O and HD^{16}O (Altwegg et al. 2015) and to separate N_2 from CO (Rubin et al. 2015).

2. DFMS measurements in the coma of 67P/Churyumov-Gerasimenko

Two different periods were chosen, 19–24 October 2014 at a distance of ~ 10 km from the center of the comet during almost circular orbits (3.15 au from the Sun) and 21–25 December 2015 (1.95 au from the Sun). The first period contains measurements covering both cometary hemispheres (-47° to 50° sub-spacecraft latitude, Jorda et al. 2015) and several cometary rotations, while the latter consists of measurement of positive latitudes (34° – 70°) covering several cometary rotations at a distance of 75–92 km from the center of the nucleus. The first measurement period was chosen because of the close distance to the comet and therefore high abundances of the species in the coma (especially $^{13}\text{C}^{16}\text{O}_2$ and $^{12}\text{C}^{18}\text{O}^{16}\text{O}$). The second period was specifically chosen because of its relatively low abundance in CO_2 , resulting in measurements for m/z 44–46 u/e on the same gain step.

3. Data treatment

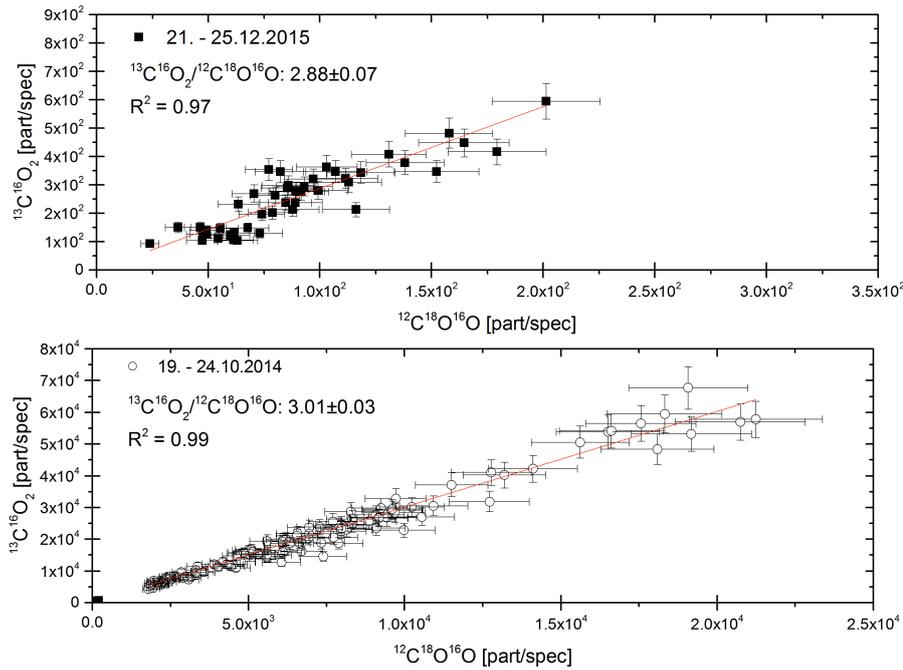
To deduce the isotopic ratios of $^{16}\text{O}/^{18}\text{O}$ and $^{12}\text{C}/^{13}\text{C}$ in CO_2 , measurement of mass per charge 44 u/e, 45 u/e, and 46 u/e are recorded. These masses are measured one after the other during normal measurement scans. The time difference in recording time between the first and the last mass is less than two minutes; therefore the compositional changes in the coma (cf. rotation period ~ 12 h by Mottola et al. 2014) between measurements of these three masses are assumed to be negligible. Each individual spectrum was corrected for the individual pixel gain of the MCP for the period in question.

The peak of mass per charge 44 u/e is dominated by the signal of CO_2 . The signal for CS is two orders of magnitude smaller than the signal for CO_2 (see Fig. 2) but still taken into account when calculating the signal for CO_2 . A separation of several peaks is therefore not needed and the signal for CO_2 is taken as the sum of the signal under the whole peak. The signal on m/z 45 u/e and 46 u/e were fitted using the above-mentioned method. For the signal for $^{13}\text{C}^{16}\text{O}_2$, $^{12}\text{C}^{18}\text{O}^{16}\text{O}$ and CO_2 , the sum over the fitted peak is used.

Background contamination of the signal due to spacecraft outgassing (Schlappi et al. 2010) was subtracted from the measured peak signals. For the October 2014 data set, the spacecraft background was determined from the signal of 1 August 2014, when the spacecraft was still far away from the comet and the coma signal was not yet detected. The subtracted spacecraft background for the December 2015 period was determined as the lowest signal for CO_2 during this time period ($2\times$ smaller than

Table 2. Results of the three different analysis methods to determine the isotopologues of CO₂ in the coma of 67P/Churyumov-Gerasimenko measured with ROSINA/DFMS.

$^{13}\text{C}^{16}\text{O}_2/^{12}\text{C}^{18}\text{O}^{16}\text{O}$	Sum of the signal	Co-added spectra	Linear fit
19.–24. October 2014	2.95 ± 0.17		3.01 ± 0.03
21.–25. December 2015	2.85 ± 0.17	3.00 ± 0.21	2.88 ± 0.07
$^{12}\text{C}^{18}\text{O}^{16}\text{O}/^{12}\text{C}^{16}\text{O}_2$	Sum of the signal	Co-added spectra	Linear fit
21.–25. December 2015	$(4.13 \pm 0.22) \times 10^{-3}$	$(3.93 \pm 0.28) \times 10^{-3}$	$(4.09 \pm 0.09) \times 10^{-3}$
$^{13}\text{C}^{16}\text{O}_2/^{12}\text{C}^{16}\text{O}_2$	Sum of the signal	Co-added spectra	Linear fit
21.–25. December 2015	$(1.19 \pm 0.07) \times 10^{-2}$	$(1.18 \pm 0.08) \times 10^{-2}$	$(1.19 \pm 0.02) \times 10^{-2}$


Fig. 3. Signal for $^{13}\text{C}^{16}\text{O}_2$ versus $^{12}\text{C}^{18}\text{O}^{16}\text{O}$ in the coma of 67P. The *top panel* black squares show the correlation of the signal for 21–25 December 2015. The *bottom panel* shows the data set of October 2014 (circles) and the measurements of December 2015 (the black squares in the left corner). Error bars show uncertainties of the single measurements. The red line is the least square linear fit.

the signal for background determined in 2014) and terrestrial abundances for the signal of the other isotopes were determined. The signal for CO₂ is two orders of magnitude higher than the signal for the other (rarer) isotopologues, resulting mostly in measurements on different gain steps and therefore an additional correction has to be taken into account, which adds to the total uncertainty.

To reduce the influence of the gain step, the ratio for m/z 46/45 is used. The signals for $^{12}\text{C}^{18}\text{O}^{16}\text{O}$ and $^{13}\text{C}^{16}\text{O}_2$ are similar in abundance (within a factor of ~ 3) and therefore always measured on the same gain step. The remaining correction is then only due to the individual pixel gain, and since the difference in pixel is never more than three pixels this choice of isotopologues results in a very accurate measurement of the combined isotopic ratio of $^{13}\text{C}^{16}\text{O}_2/^{12}\text{C}^{18}\text{O}^{16}\text{O}$.

4. Results

We divide the measured isotopic ratios for CO₂ into three sections: $^{13}\text{C}^{16}\text{O}_2/^{12}\text{C}^{18}\text{O}^{16}\text{O}$, $^{16}\text{O}/^{18}\text{O}$, and $^{12}\text{C}/^{13}\text{C}$. For the data from the time period in December 2015 (the time where the gain step was similar for each consecutive mass spectrum at m/z 44, 45, 46), three different analysis methods were chosen to calculate the isotopic ratios and therefore provide comparison: a) each individual spectrum was fitted to separate the signals of

the different isotopologues in overlapping peaks. Then the total signal for the individual isotopologue was derived by adding up the contribution from all spectra and then divided by the total signal of the other isotopologue to determine the ratio; b) all spectra were co-added based on the mass scale, resulting in a single spectrum for each mass. The signal for the individual isotopologue is determined by fitting the overlapping peaks resulting in signals for each isotopologue; c) each individual spectrum was fitted to separate the isotopologue signal from overlapping peaks. From each triplet of back-to-back measurements of mass per charge 44 u/e, 45 u/e, and 46 u/e isotopic ratios were derived to investigate possible correlations. The slope of the linear fitted curve corresponds then to the measured isotopic ratio. The linear fit takes into account the individual uncertainties of the measurements resulting in a least squares fit.

The total uncertainty of the sum of the signal is composed of the ion statistical counting uncertainty and the pixel gain uncertainty of 5%. The total uncertainty of the co-added spectra takes into account the peak fitting uncertainty (7%) and the ion statistical uncertainty. The uncertainty of the linear fit of the correlation is the standard error of the fit.

Due to stochastic effects, the oxygen isotope has a two times higher probability of being found in the signal for $^{12}\text{C}^{18}\text{O}^{16}\text{O}$, therefore a factor of 2 has to be taken into account (compare results for linear fit in Figs. 4 to 5). Table 2 shows an overview

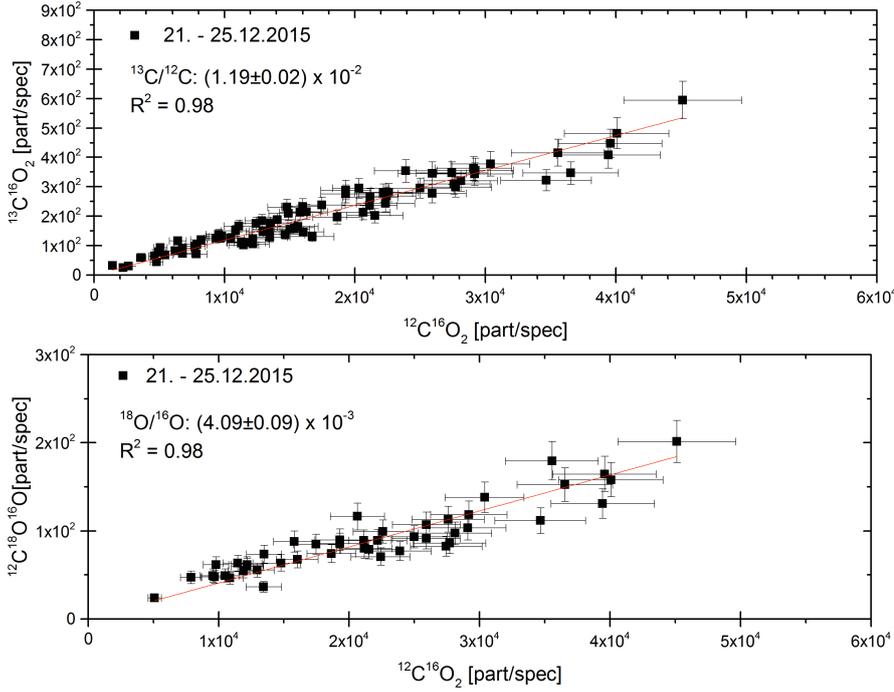


Fig. 4. Measured signal on the detector for $^{13}\text{C}^{16}\text{O}_2$ and $^{12}\text{C}^{18}\text{O}^{16}\text{O}$ versus $^{12}\text{C}^{16}\text{O}_2$. The red line represents the resulting linear fit.

of the measured ratios in CO_2 for the sum of the signal, the co-added spectra, and the linear fit. The results of the individual analysis methods are all overlapping within uncertainties and show that the results are independent of the analysis method.

$^{13}\text{C}^{16}\text{O}_2/^{12}\text{C}^{18}\text{O}^{16}\text{O}$ in CO_2 : in addition to the December 2015 data, the measurements for $^{12}\text{C}^{18}\text{O}^{16}\text{O}/^{13}\text{C}^{16}\text{O}_2$ in CO_2 for the time period of October 2014 are all measured on the same gain step, therefore a total of five measurements (two sum of the signal, two linear fit, and one co-added spectra) are provided. The least square linear fit for the time period of October 2014 and December 2015 is shown in Fig. 3. The measurements of the two time periods show a good correlation ($R^2 > 0.97$). If there were contamination of the $^{12}\text{C}^{18}\text{O}^{16}\text{O}$ signal by NO_2 , the correlation is expected to show this due to the compositional changes in the coma between the northern and southern hemisphere (Le Roy et al. 2015). The good correlation therefore shows that either there is no or negligible contamination of the signal with NO_2 or the signal would be perfectly correlated with CO_2 . The results of the different analysis methods for the two time periods are overlapping within uncertainties. The uncertainty of the average takes into account the individual uncertainties as well as the standard uncertainty of the average. The average of these three methods for two time periods and resulting uncertainty for $^{13}\text{C}^{16}\text{O}_2/^{12}\text{C}^{18}\text{O}^{16}\text{O}$ in the coma of 67P is 2.93 ± 0.04 .

$^{16}\text{O}/^{18}\text{O}$: the results of the three different analysis methods (least square linear fit in Fig. 4) are overlapping within uncertainties. The uncertainty of the average takes into account the individual uncertainties as well as the standard uncertainty of the average. The average of these three methods and resulting uncertainty for $^{18}\text{O}/^{16}\text{O}$ in CO_2 in 67P's coma is 494 ± 8 ($(2.03 \pm 0.03) \times 10^{-3}$). The average itself shows slight enrichment in ^{18}O , but is within uncertainties still consistent with the terrestrial value by Meija et al. (2016) (499) and also in agreement with earlier measurements in H_2O by ROSINA/DFMS 556 ± 62 Altwegg et al. (2015).

$^{12}\text{C}/^{13}\text{C}$: the ratios of the three different analysis methods (least squares linear fit in Fig. 4) are overlapping within

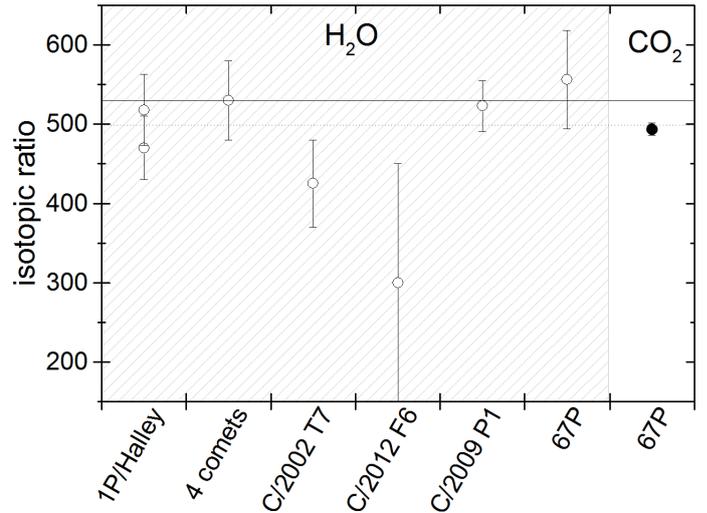


Fig. 5. Measurements of $^{16}\text{O}/^{18}\text{O}$ in water and carbon dioxide for several comets (for references see Table 3). The measurement of $^{16}\text{O}/^{18}\text{O}$ in CO_2 shows a significantly smaller error bar than earlier measurements (black circle) and is within uncertainties compatible with a terrestrial abundance by Meija et al. (2016) (dotted line $^{16}\text{O}/^{18}\text{O}$) but not with a solar abundance by McKeegan et al. (2011) of 530 (solid line).

uncertainties. The uncertainty of the average takes into account the individual uncertainties as well as the standard uncertainty of the average. The average of these three methods and resulting uncertainty for $^{12}\text{C}/^{13}\text{C}$ in CO_2 in 67P's coma is 84 ± 4 ($(1.19 \pm 0.05) \times 10^{-2}$). This ratio is not in agreement with the solar wind abundance by Hashizume et al. (2004) nor with the terrestrial standard provided by Zhang et al. (1990), although the latter is almost compatible. The $^{12}\text{C}/^{13}\text{C}$ for CO_2 in 67P's coma by ROSINA/DFMS shows an enrichment in ^{13}C of 6% compared to terrestrial abundances by Meija et al. (2016) and a 14% enrichment compared to solar abundances, which is not compatible even within uncertainties (1σ). Other measurements with

Table 3. Isotopic abundances in comets adopted from Jehin et al. (2009), Woods (2009), and Bockelée-Morvan et al. (2015).

Isotopic ratio	Species	Value	Comet	Reference
¹² C/ ¹³ C	C ₂	60 ± 15	West 1976 VI	Lambert & Danks (1983)
		93 ± 10	4 comets	Wyckoff et al. (2000)
		85 ± 20	C/2002 T7	Rousselot et al. (2012)
		80 ± 20	C/2001 Q4	Rousselot et al. (2012)
		94 ± 33	C/2012 S1	Shinnaka et al. (2014)
	CN	65 ± 9	1P/Halley	Wyckoff et al. (1989)
		89 ± 17	1P/Halley	Jaworski & Tatum (1991)
		95 ± 12	1P/Halley	Kleine et al. (1995)
		90 ± 15	C/1995 O1	Lis et al. (1997)
		165 ± 40	C/1995 O1	Arpigny et al. (2003)
		90 ± 10	C/1990 K1 Levy	Wyckoff et al. (2000)
		85 ± 20	C/1989 X1 Austin	Wyckoff et al. (2000)
		93 ± 20	C/1989 XIX Okazaki-Levy-Rudenko	Wyckoff et al. (2000)
		90 ± 10	122P/1995 S1 de Vico	Jehin et al. (2004)
		90 ± 10	88P/1981 Q1 Howell	Hutsemékers et al. (2005)
	HCN	95 ± 15	9P/ Temple 1	Jehin et al. (2006)
		90 ± 20	17P/Holmes	Bockelée-Morvan et al. (2008)
		91 ± 4	21 comets	Manfroid et al. (2009)
		95 ± 15	103P/Hartley 2	Jehin et al. (2011)
		95 ± 25	C/2012 F6	Decock et al. (2014)
111 ± 12		C/1995 O1	Jewitt et al. (1997)	
94 ± 8		C/1995 O1	Bockelée-Morvan et al. (2008)	
C ₂ H ₄	114 ± 26	17P/Holmes	Bockelée-Morvan et al. (2008)	
	109 ± 14	C/2014 Q2	Biver et al. (2016)	
	124 ± 64	C/2012 F6	Biver et al. (2016)	
	83 ± 12	67P	Rubin et al. (2017)	
	84 ± 9	67P	Rubin et al. (2017)	
¹⁶ O/ ¹⁸ O	CO	86 ± 9	67P	Rubin et al. (2017)
		84 ± 4	67P	This work
		518 ± 45	1P/Halley	Balsiger et al. (1995)
		470 ± 40	1P/Halley	Eberhardt et al. (1995)
		530 ± 50	4 comets	Biver et al. (2007)
CO ₂	425 ± 55	C/2002 T7	Hutsemekers et al. (2008)	
	300 ± 150	C/2012 F6	Decock et al. (2014)	
	523 ± 32	C/2009 P1	Bockelée-Morvan et al. (2012)	
	556 ± 62	67P	Altwegg et al. (2015)	
	494 ± 8	67P	This work	
	2.93 ± 0.04	67P	This work	

ROSINA/DFMS in different molecules, namely C₂H₄, C₂H₅, and CO show a similar behavior for ¹²C/¹³C (including fragments of CO₂) (Rubin et al. 2017). However, these measurements provide a substantially larger error bar than presented for the ¹²C/¹³C in CO₂ in 67P's coma as their investigated period was specifically chosen to analyze silicon isotopes and shows large interferences between ¹³CO and HCO, resulting in larger uncertainties.

5. Discussion

The measurements presented show a slight enrichment in ¹³C in the ¹²C/¹³C ratio which is not in agreement with measurements by remote sensing of carbon isotopes in other molecules (see Table 3 and Fig. 1). Contrary to the DFMS results at 67P, two measurements of HCN show even a depletion of ¹³C compared to the solar system abundance. Measurements of the carbon isotope in HCN are not feasible for DFMS due to overlapping peaks with DCN and CH₂N. An overview on

feasible isotopic measurements by ROSINA/DFMS is given by Bockelée-Morvan et al. (2015). However, in most cases error limits are too large to make these differences significant. Our ratio is, on the other hand, compatible with measurements in CO (including 1/3 fragments of CO₂) and C₂H₄ and C₂H₅ (which both are most probably fragments of ethane) in 67P's coma by ROSINA/DFMS (Rubin et al. 2015).

For the ¹⁶O/¹⁸O ratio, the DFMS value is compatible with a terrestrial value within the error limit of 2%, the most precise value so far for the ¹⁶O/¹⁸O ratio in comets. If there is a depletion or enrichment of the heavy oxygen isotope in CO₂ it has to be small. The measurements of the combined carbon and oxygen isotopes (¹³C¹⁶O₂ and ¹²C¹⁸O¹⁶O) are, for the two time ranges, compatible within uncertainties. However, the focus for the selection of those data sets was based on minimizing instrumental uncertainties to detect time dependent variations or heterogeneity in the coma of the comet and would require a careful analysis of the total data available, which is not the scope of this paper.

6. Conclusion

The isotopic measurements for CO₂ are the first ever presented in situ measurements of ¹²C/¹³C and ¹⁶O/¹⁸O in CO₂ for a cometary coma. The derived ¹⁶O/¹⁸O ratio of 494 ± 8 is compatible with a terrestrial abundance of 499 by Meija et al. (2016) but not with a solar abundance of 530 by McKeegan et al. (2011). The ¹²C/¹³C ratio of 84 ± 4 (1.19 ± 0.05)% is 6% (1σ) lower than the terrestrial value of 1.12% and 14% higher than solar wind abundance, although given the slight difference between the upper limit of the measured value and the terrestrial standard it could be considered compatible with a terrestrial origin. It is in very good agreement with the value derived in 67P for ¹²C/¹³C in CO that contains the signal of the CO₂ fragment and the parent (Rubin et al. 2017). The model by Woods & Willacy (2009) predicts a slight carbon fractionation in the midplane at distances inside of 25AU as well as for material vertically offset for CO and CO₂. The model prediction of the ¹²C/¹³C ratio in HCN for comets is not compatible with most measurements in comets provided in Table 3. Several possible explanations are provided by Woods & Willacy (2009). One explanation is that the protosolar disk was heated to a temperature high enough so that carbon isotopic ratios are mixed to the extent that carbon isotopic abundances in all molecules is similar. This explanation can probably be ruled out because measurements of D/H in water taken at the same time show a very high ratio for 67P (Altwegg et al. 2015) compared to other measurements in comets or terrestrial values. With the same argument, the possibility of relying on mixing material up to the surface layers of the disk (dissociating the molecules), which then leads to a reset of the carbon isotope ratio, can also be ruled out. That is, unless there is a difference of volatility between CO and CO₂ relative to water because water is frozen out on grains. Then it can be argued that grains did not mix, whereas gas did. In addition, the third explanation of a passing shock affecting the protosolar disk proposed by Woods & Willacy (2009) does not seem plausible for the solar system as again all material would have passed through this shock and no fractionation of isotopes would be seen in cometary volatiles. The small fractionation of ¹²C/¹³C in CO and CO₂, together with some possible fractionation in HCN (reference in Table 3), is probably compatible with the ¹²C/¹³C ratio inherited from the original protosolar cloud. That means that according to the model by Woods & Willacy (2009), 67P formed at distances further away from the Sun than 25 AU. This formation location is then in agreement with D/H in water, D₂O, and HDS (Altwegg et al. 2017) which all point to a heritage from the native cloud in the protosolar nebula.

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paper are present in the paper. Additional data related to this paper may be requested from the authors. The authors would like to thank the referees for constructive criticism, comments, and input, which helped to improve the clarity and quality of the paper.

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